Innovations in Tunable Optical Switches: Permittivity Control in Metals, Semiconductors and Dielectrics

Soham Saha

Ph.D., Electrical Engineering,

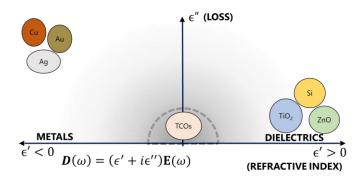
Purdue University (2021)



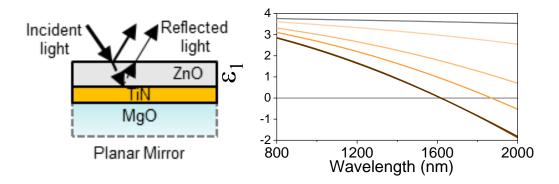




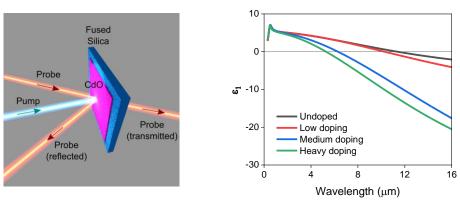
Outline



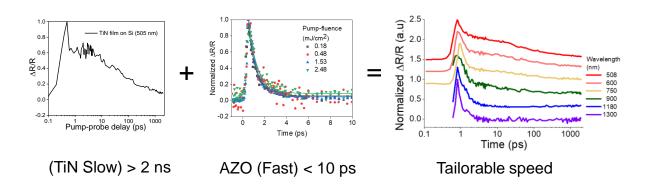
Permittivity and the Epsilon near Zero Point: How to **tune** or **tailor** them



How much can we **tune** the permittivity of zinc oxide **without** adding dopants?

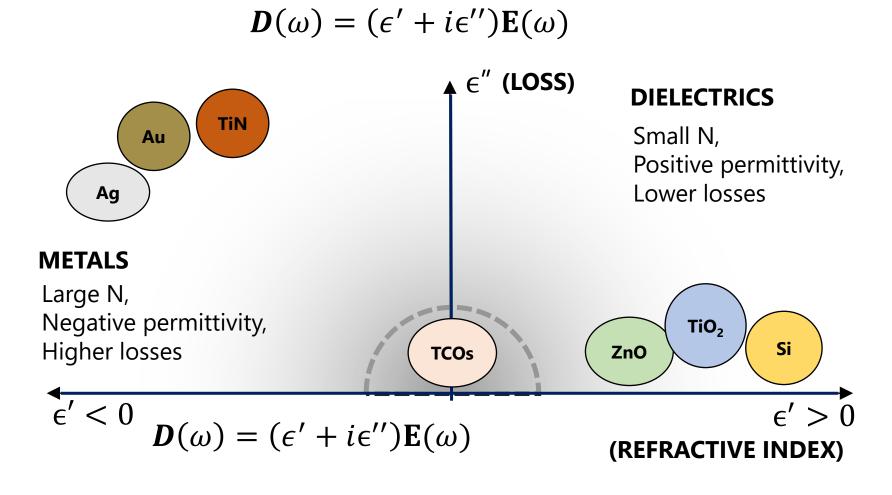


How much can we **tailor** the permittivity and relaxation-time of cadmium oxide by adding dopants?



Can we **control** the overall **response time** of a metasurface?

Effect of Free-Electrons on Permittivity



Drude Lorentz Model

$$w_p^2 = \frac{Ne^2}{\epsilon_0 m^*}$$

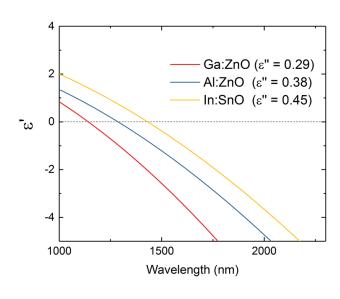
$$\epsilon' = \epsilon_{\rm B} - \frac{w_{\rm p}^2}{w^2 + \Gamma_d^2}$$

$$\epsilon'' = \frac{w_p^2 \Gamma_d}{w^3 + \Gamma_d^2 w}$$

Epsilon-Near-Zero (ENZ) Effects

ENZ regime is the wavelength range where the permittivity changes sign Happens near telecom freq. in TCOs

Reflectance modulation enhanced near ENZ



$$n = \sqrt{\epsilon}$$

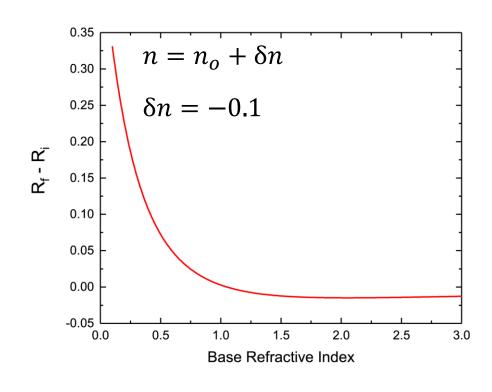
$$\delta n = \delta \epsilon / (2\sqrt{\epsilon})$$

$$\delta R = \frac{1}{2\sqrt{\epsilon}} \frac{dR}{dn} \delta \epsilon$$

$$\epsilon \to 0$$
 $\delta R \to \infty$

Large Field Enhancements $\varepsilon_1 \mathbf{E}_1 = \varepsilon_2 \mathbf{E}_2$ **Slow light effects** => Enhanced light-matter interaction

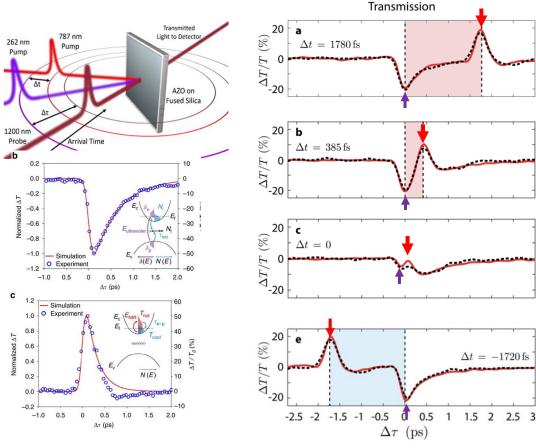
$$v_g = c\sqrt{\varepsilon'(\omega)} \left(1 + \frac{2}{\pi} \int_0^\infty \frac{\varepsilon''(\omega_1)}{(\omega_1^2 - \omega^2)^2} \omega_1^3 d\omega_1 \right)^{-1}.$$
 Khurgin et al. *Optica* (2020)



Kinsey et al. *Optica* (2015)

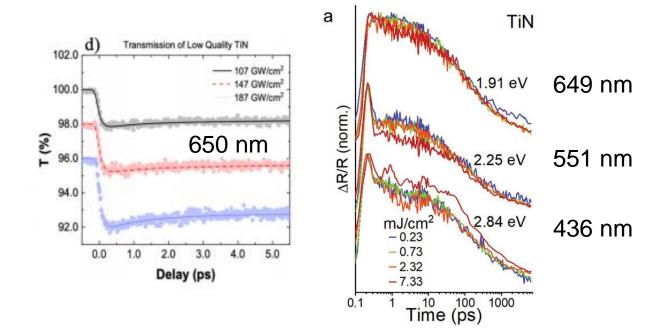
Application of Large ENZ-Enhanced Reflectance Modulation

All optical addition of signals at femtosecond timescales



Clerici et al. Nat. Comm. (2016)

Extraction of hidden dynamics of hot-electrons



H. George et al. *Opt. Mat. Exp.* (2019)

B. Diroll et al. *Adv. Opt. Mat.* (2020)

Nonlinearities Enhanced by ENZ

Large optical nonlinearity of indium tin oxide in its epsilon-near-zero region

M. Zahirul Alam¹, Israel De Leon^{1,3,*}, Robert W. Boyd^{1,2}

+ See all authors and affiliations

Science 13 May 2016: Vol. 352, Issue 6287, pp. 795-797 DOI: 10.1126/science.aae0330



Science Vol 352, Issue 6287 13 May 2016 Table of Contents Print Table of Contents Advertising (PDF) Classified (PDF) Masthead (PDF)

Second Harmonic Generation from Phononic Epsilon-Near-Zero Berreman Modes in Ultrathin Polar Crystal Films

Nikolai Christian Passler, I. Razdolski, D. Scott Katzer, D. F. Storm, Joshua D. Caldwell, Martin Wolf, and Alexander Paarmann*

Cite this: ACS Photonics 2019, 6, 6, 1365-1371 Publication Date: June 3, 2019 >

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Broadband frequency translation through time refraction in an epsilon-near-zero material

Yiyu Zhou I, M. Zahirul Alam, Mohammad Karimi, Jeremy Upham, Orad Reshef, Cong Liu, Alan E. Willner & Robert W. Boyd

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Field-Effect Tunable and Broadband Epsilon-Near-Zero Perfect Absorbers with Deep Subwavelength Thickness

Aleksei Anopchenko*, Long Tao, Catherine Arndt, and Ho Wai Howard Lee*

♥ Cite this: ACS Photonics 2018, 5, 7, 2631-2637 Publication Date: April 23, 2018 > https://doi.org/10.1021/acsphotonics.7b01373 Copyright © 2018 American Chemical Society

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Negative Refraction in Time-Varying Strongly Coupled Plasmonic-Antenna-Epsilon-Near-Zero Systems

V. Bruno, C. DeVault, S. Vezzoli, Z. Kudyshev, T. Hug, S. Mignuzzi, A. Jacassi, S. Saha, Y. D. Shah, S. A. Maier, D. R. S. Cumming, A. Boltasseva, M. Ferrera, M. Clerici, D. Faccio, R. Sapienza, and V. M. Shalaev Phys. Rev. Lett. 124, 043902 - Published 30 January 2020

Physics See Synopsis: Plasmonic Metamaterials Bend Light Backwards

Published: 01 May 2017

Femtosecond optical polarization switching using a cadmium oxide-based perfect absorber

Yuanmu Yang M. Kyle Kelley, Edward Sachet, Salvatore Campione, Ting S. Luk, Jon-Paul Maria, Michael B. Sinclair & Igal Brener □

Nature Photonics 11, 390–395(2017) | Cite this article

2406 Accesses | 108 Citations | 60 Altmetric | Metrics

High-harmonic generation from an epsilon-nearzero material

Yuanmu Yang 🖂, Jian Lu, Alejandro Manjavacas, Ting S. Luk, Hanzhe Liu, Kyle Kelley, Jon-Paul Maria, Evan L. Runnerstrom, Michael B. Sinclair, Shambhu Ghimire & Igal Brener

Nature Physics 15, 1022-1026(2019) Cite this article 6416 Accesses 28 Citations 15 Altmetric Metrics

Review Article | Published: 26 September 2019

Near-zero-index materials for photonics

Nathaniel Kinsey ☑, Clayton DeVault, Alexandra Boltasseva & Vladimir M. Shalaev ☑

Nature Reviews Materials 4, 742-760(2019) Cite this article

4468 Accesses 41 Citations 14 Altmetric Metrics

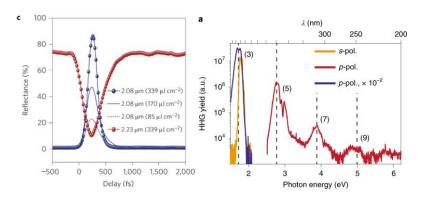
Can we **extend Epsilon Near Zero** effects to any wavelength of our choice?

How much can we **control** the permittivity in **real time**?

Can we **control** the **response-time** of an all-optical switch?

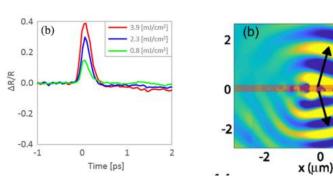
Materials Explored in This Work

Cadmium Oxide (CdO)



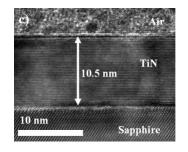
- Femtosecond switching: Nat. Phot. (2017)
- High harmonic generation: Nat. Phys (2019)
- Low-loss metal for MID IR Plasmonics
- High mobility for electroabsorption modulators
- **ENZ Applications**: ultrafast switches, high harmonic Generation

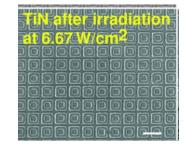
Zinc Oxide (ZnO) and Al:ZnO



- Kinsey et. al. Optica (2015)
- V. Bruno et. al. PRL. (2020)
- ZnO: laser tolerant, used in optoelectronic applications from single photon emitters to solar cells
- AZO Demonstrated in various ENZ applications like reflectance enhancement, negative refraction, etc.

Titanium Nitride (TiN)





Shah et. al. AOM. (2017)

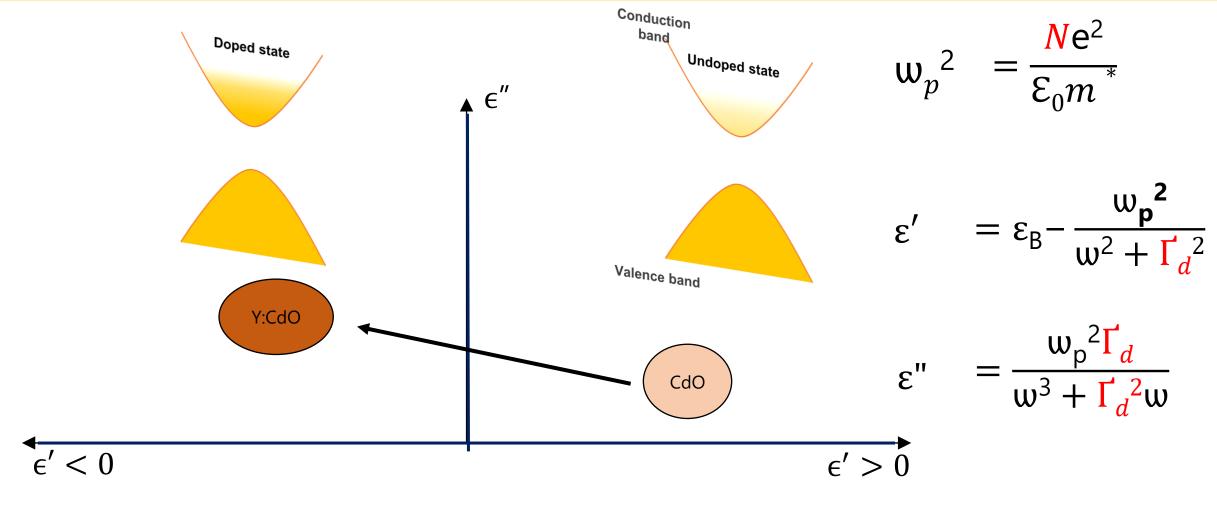
Guler et. al. Adv. Mat. (2014)

- Gold-like metallic properties
- Large laser tolerance
- Interesting dynamics with a nanosecond response time



How much can we tailor the **permittivity** and **relaxation time** of cadmium oxide **by adding dopants**?

How do We Tailor the Permittivity?

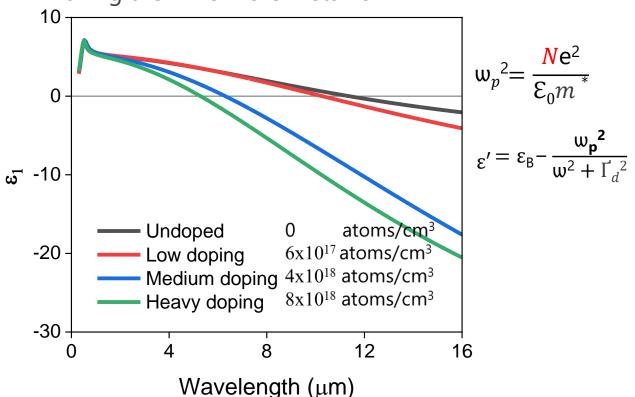


$$Y_2O_3 \xrightarrow{CdO} 2Y_{Cd} + 2O_0^X + 2n + \frac{1}{2}O_2(g)$$

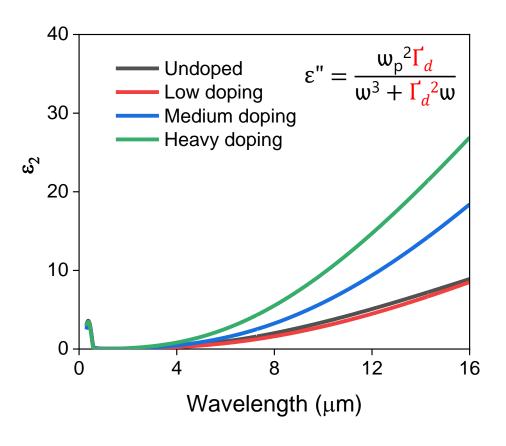
Yttrium Doping of Cadmium Oxide to Tailor the Permittivity

Material: CdO ~500 nm films grown on fused silica by co-sputtering CdO and Yttrium

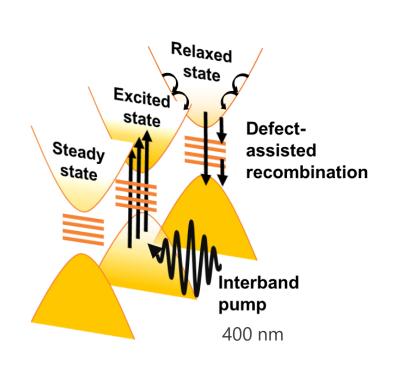
Adding more Y dopant increases electrons and blue shifts the ENZ point from 11 to 5 microns, making the films more metallic



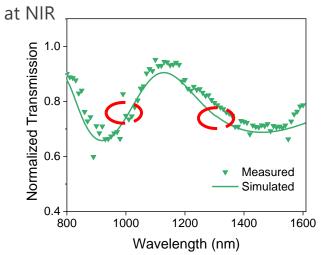
Increased electrons and defects increases the losses with doping



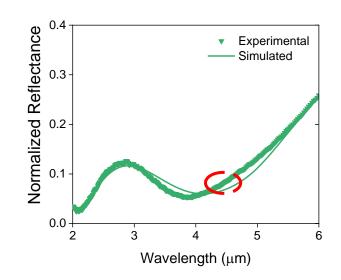
ENZ Enhanced All-optical Switching in Mid-IR



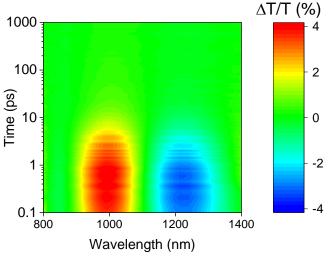
Steady state transmittance vs wavelength at NIR _____



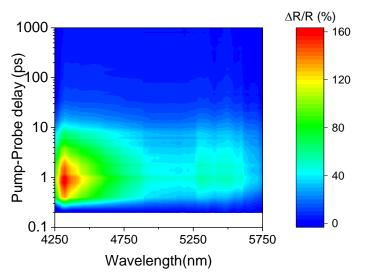
Steady state reflectance vs wavelength at MIR



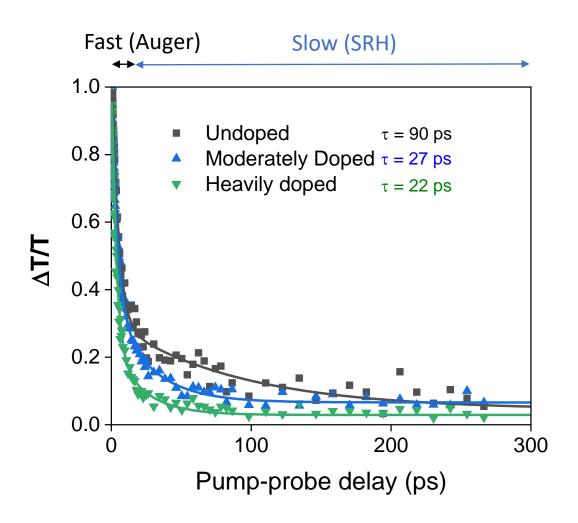
Moderate modulation near transmittance dips (4%)



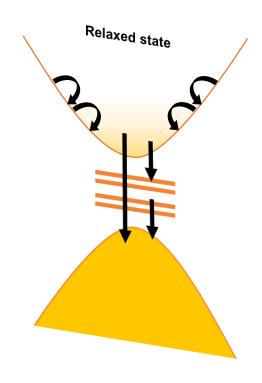
Further enhanced near the ENZ point (160%)



The Overall Response Time can be Engineered by Doping



Defect-assisted recombination

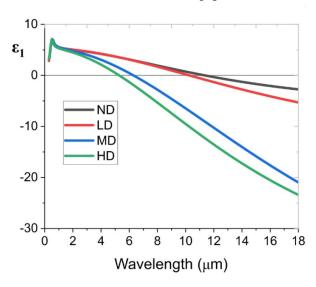


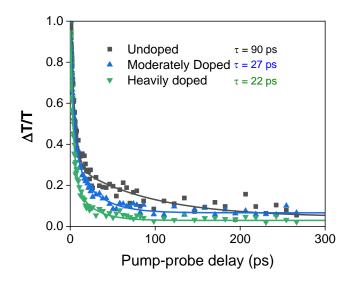
More defects, faster recombination

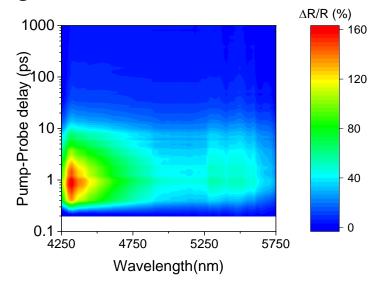
Summary

Application: All optical switching in the MIR for satellite communications

Application 2: ENZ enhanced phenomena across a broad MIR range







ENZ can be tailored from 5 to 11 microns by adding dopants

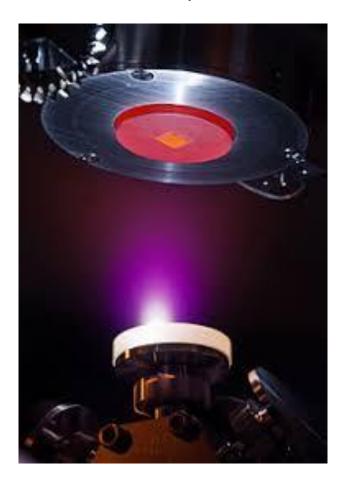
Doping reduces the relaxation time

Modulation of 160% near ENZ at 1.3 mJ/cm² pump power with 50ps relaxation time

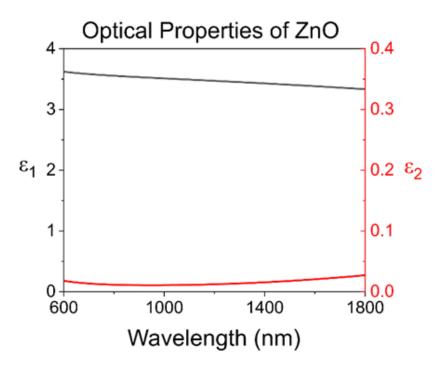
Doping introduces losses in the oxide Large on-state losses The process is irreversible How much can we tune the **permittivity** of zinc oxide in real time, without adding dopants?

Material of Choice: Zinc Oxide (Undoped)

ZnO films grown by Pulsed Laser Deposition

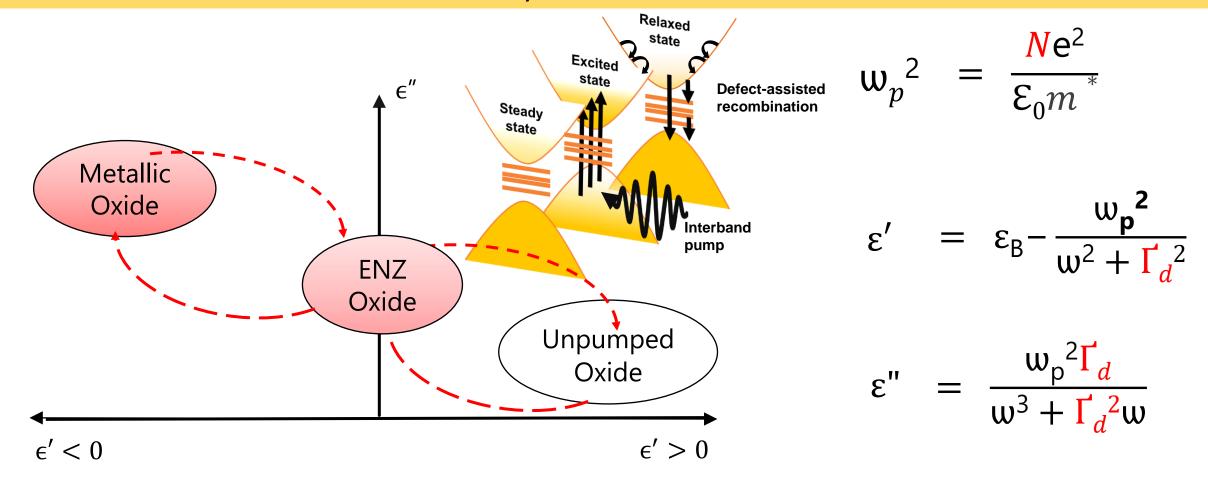


Dielectric with low optical losses



$$\varepsilon = \varepsilon_{\infty} + \frac{A_1}{E_1^2 - (\hbar\omega)^2 - iB_1\hbar\omega} - \frac{A_0}{(\hbar\omega)^2 + iB_0\hbar\omega}$$

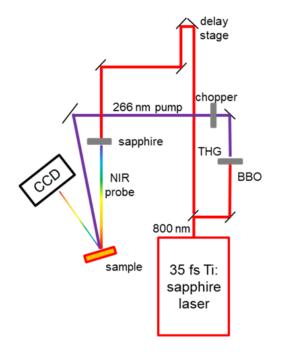
How do We Tune the Permittivity?

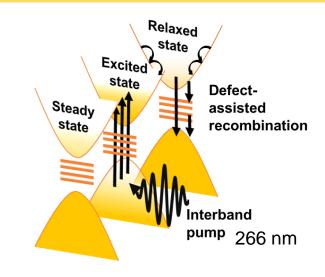


Optically pumping with an interband pump, generating electron hole pairs Tuning is a dynamic method: Instantaneous, reversible results

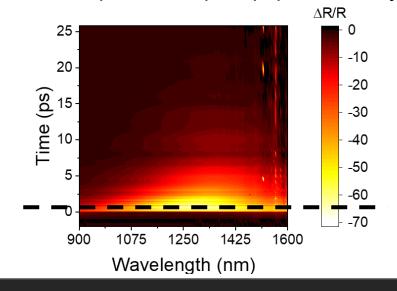
Interband Pump – Near Infrared Probe

Pump-probe schematic

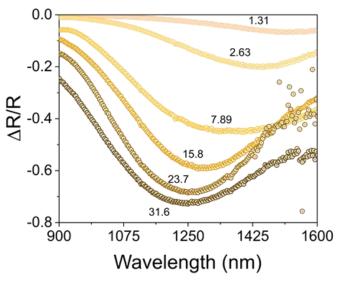




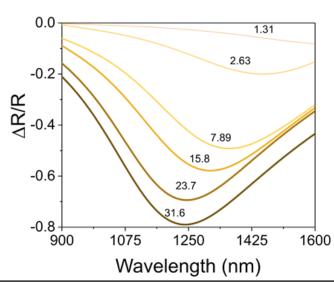
 Δ R/R captured vs pump-probe delay



Reflectance extracted

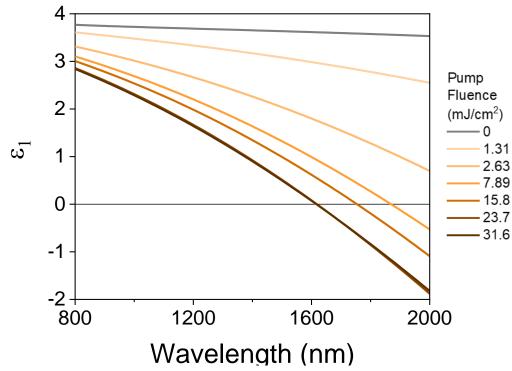


Drude Parameters extracted from fits



Unity Order Permittivity Changes at Telecom Wavelengths

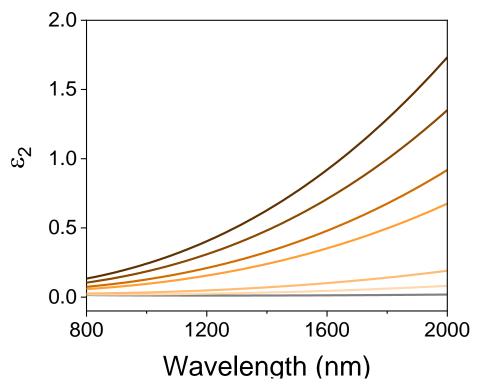
Extraordinary permittivity modulation (-4 at 1600 nm) saturates at 23 mJ/cm²



N saturates due to saturable absorption As N rises, m* rises due to non parabolic band

$$w_p^2 = \frac{Ne^2}{\varepsilon_0 m *}$$
 $\varepsilon' = \varepsilon_B - \frac{w_p^2}{w^2 + \Gamma_d^2}$

With increased power, the absorption in the films increase

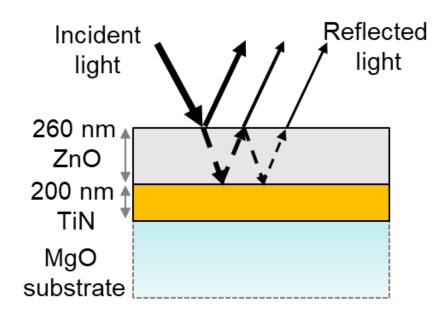


Increased free carrier concentration and heating of the lattice increases the scattering Γ

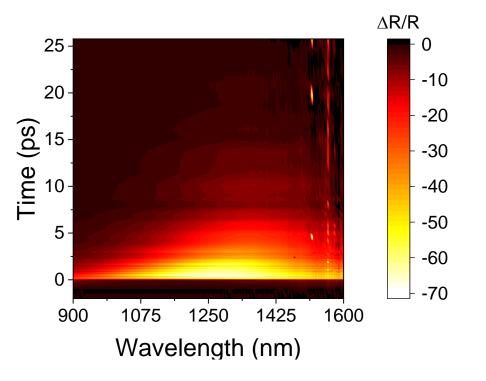
$$\varepsilon'' = \frac{\omega_p^2 \Gamma_d}{\omega^3 + \Gamma_d^2 \omega}$$

Large, Broadband Modulation in Planar Mirrors Without Lithography

Lithography-free mirror



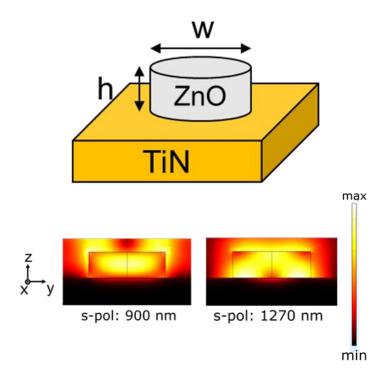
Broadband modulation of 70% at telecom Without Doping or patterning



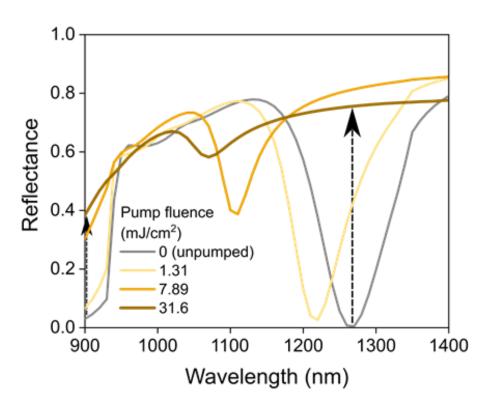
20ps response time!

Selective Enhancement at Specific Wavelengths: Design

Hybrid Plasmonic Resonators

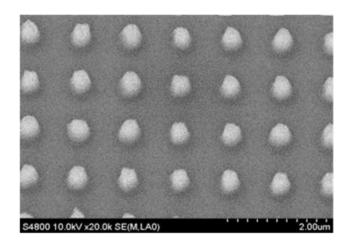


Resonance shifts under a pump

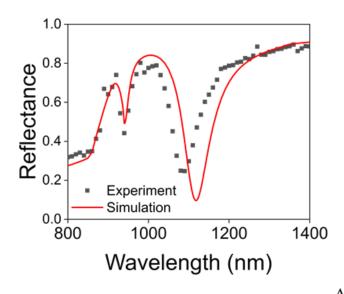


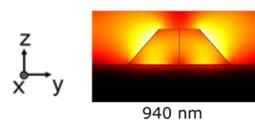
Selective Enhancement at Specific Wavelengths

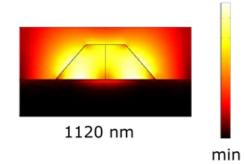




Steady-state resonance shows two dips

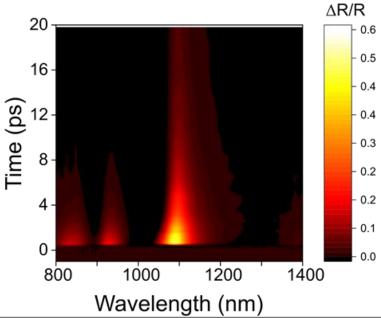






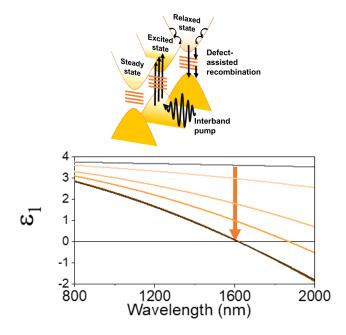
max

Maximum change of 55% at 7.8 mJ/cm² with 20 ps Response time



Summary

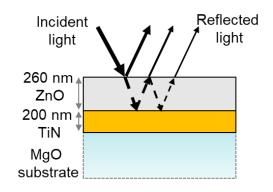
Large optical changes are demonstrated In UNDOPED ZnO by interband pumping

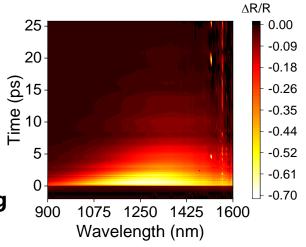


Application: All optical switching

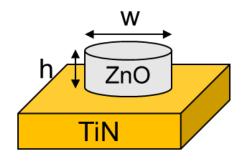
Transient permittivity data useful for designing dynamic polarization switches, HMMs, etc

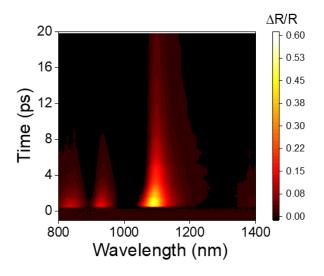
Planar metal-backed ZnO mirror Can be used to demonstrate broadband modulation with ps response time





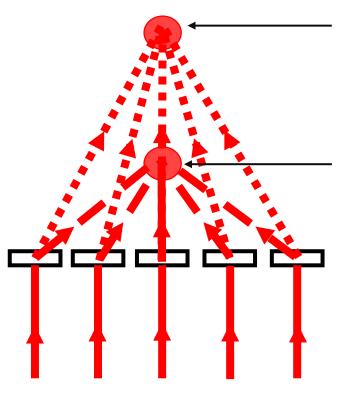
The resonances can be selectively enhanced at specific wavelengths With simple resonant structures





Optically Tunable Lens

Tunable focusing, polarization rotators, filters, etc



Focal Point at Steady State

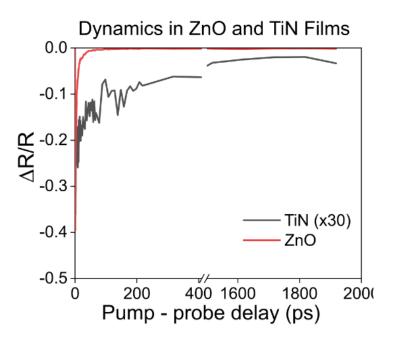
Focal Point at Pumped State

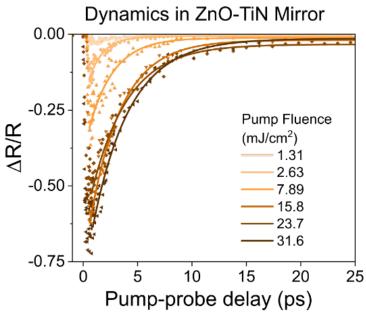
Metalens made of low-loss oxide

Impact

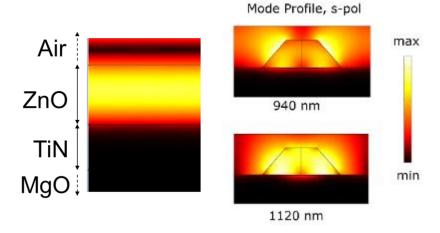
- Focusing high power beams used to burn off tissue
- All-optical scanning

The Switches are Fast, Despite the Slower Response of TiN





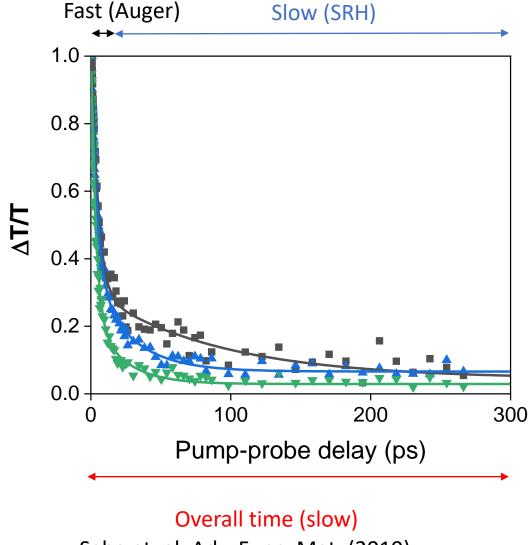
Most of the light matter interaction happens in the ZnO



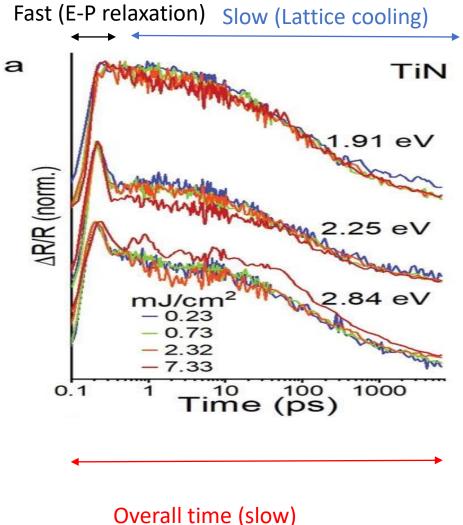
As a result, ZnO dictates the modulation dynamics, even when TiN is excited

Can we **control** the **speed** of an all-optical switch?

In Materials, Response Time is Dictated by the Slowest Dynamics



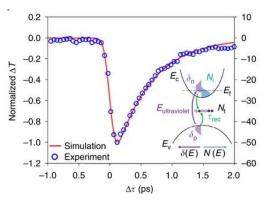
Saha et. al, Adv. Func. Mat. (2019)



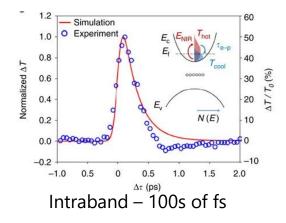
B. Diroll, S. Saha et. al, Adv. Opt. Mat. (2020)

Pump Induced Control of Switching Dynamics

Utilizing interband vs intraband Relaxation times

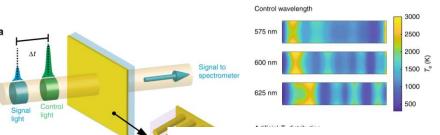


Interband – picoseconds



Clerici et al., Nat Comm. 2016

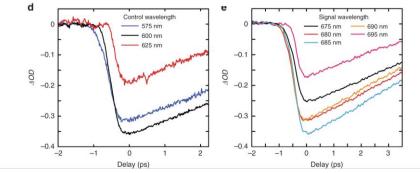
Controlling temperature distribution in nanorods controls the photon dynamics to an extent



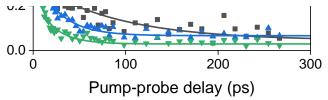
Controlling doping to control relaxation times



control the operation speed?



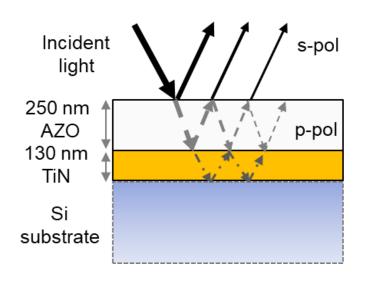
Zayats et al., Nat Comm. 2019



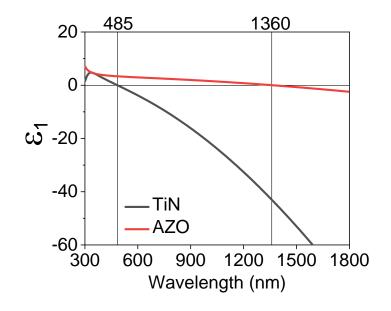
Saha et al., Adv. Opt. Mat. 2020

Berreman Metasurface Strongly Absorbs p-polarized Light Near ENZ

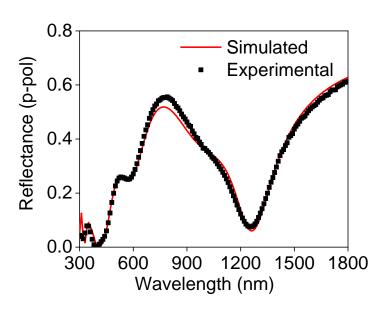
AZO film on TiN



ENZs of TiN and AZO

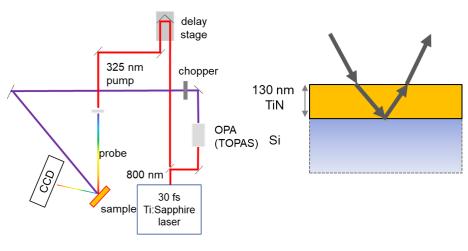


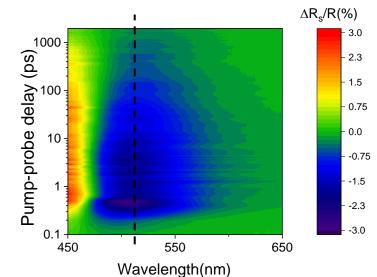
Near perfect absorption near ENZ



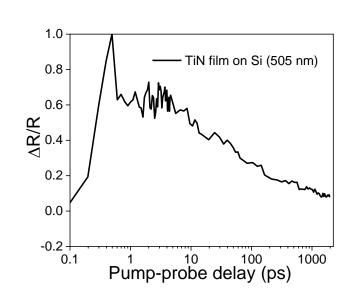
TiN is Slow (325 nm pump, Vis probes, 1.3 mJ/cm²)

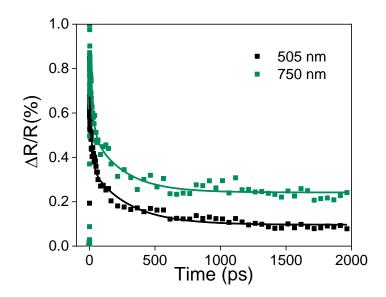






TiN Response time > 2 ns

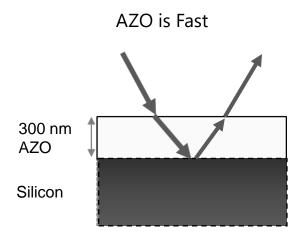


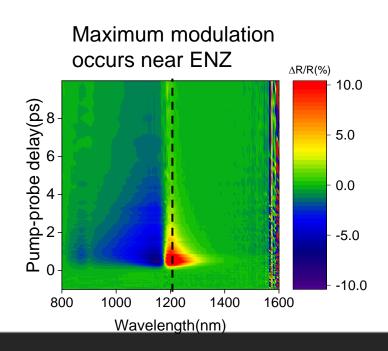


$$\left(\frac{\Delta R}{R}\right)(t) = A e^{\frac{-t}{\tau_1}} + B e^{\frac{-t}{\tau_2}} + C$$

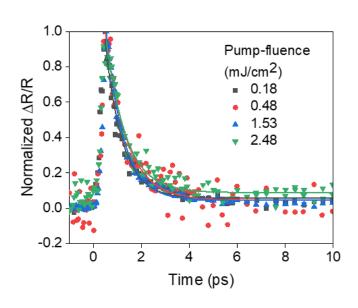
Wavelength (nm)	A	В	С	$ au_1$ (ps)	$ au_2$ (ps)
505	0.33	0.26	0.096	17.3	287
750	0.29	0.28	0.24	24.3	242

AZO is Fast (325 nm pump, NIR probes, 1.3 mJ/cm²)





Sub 10 ps response time

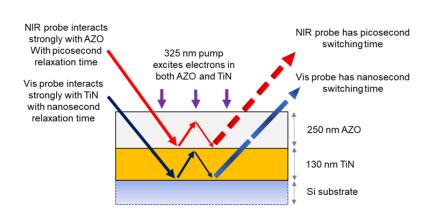


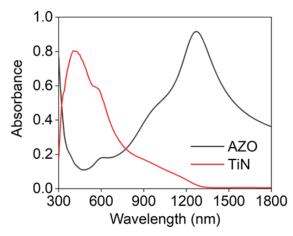
$$\left(\frac{\Delta R}{R}\right)(t) = D e^{\frac{-t}{\tau}} + E$$

Pump Fluence (mJ/cm²)	D	Е	au (ps)
0.18	1.76	0.056	0.715
0.49	1.69	0.046	0.873
1.53	1.98	0.045	0.723
2.48	1.65	0.088	0.834

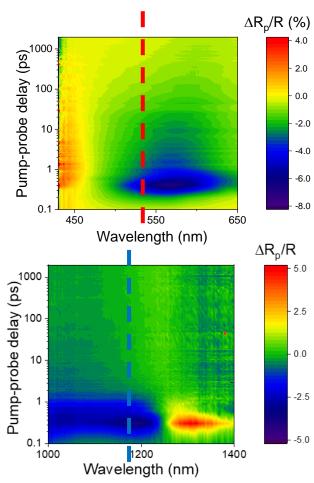
Metasurface Can Operate Both Fast and Slow

In the metasurface, light is strongly absorbed by the individual material close to the ENZ point

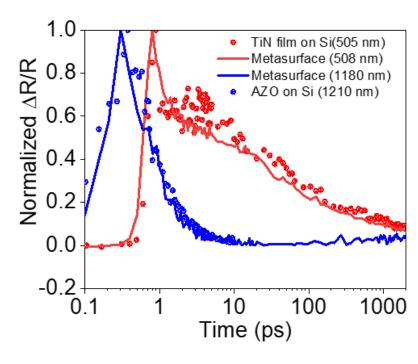




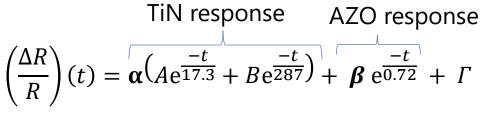
The metasurface reacts slowly near 500 nm, and fast near 1300 nm.



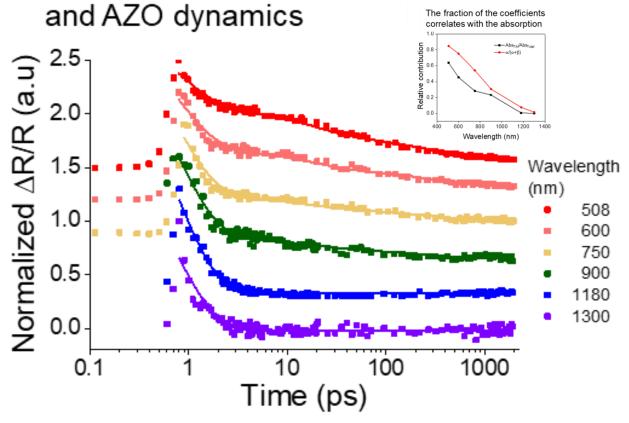
The metasurface switching responses (solid lines) near ENZ match well with that of the individual films (dots)

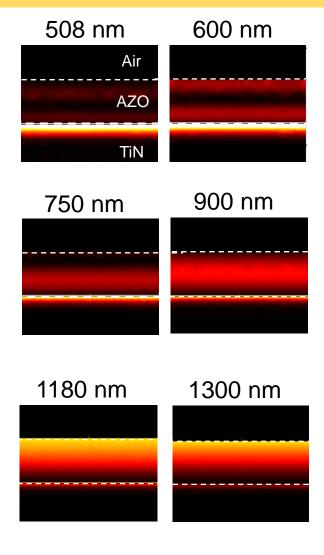


Designer Dynamics Determined by the Light Matter Interaction



The decay as a weighted sum of the TiN





Final Summary

Understanding tuning and tailoring can help us understand the steady state and temporal dynamics of TCOs A deeper understanding of the materials will help us develop new devices

Y:CdO

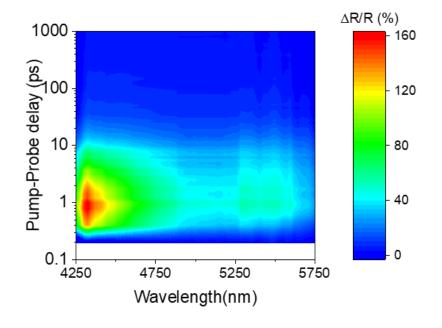
- ENZ Tailoring with different doping
- Controlling **decay time** with dopants
- Large modulation in the MIR

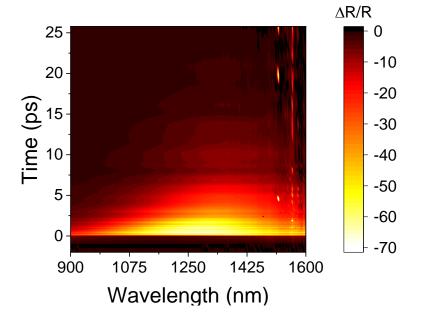
ZnO

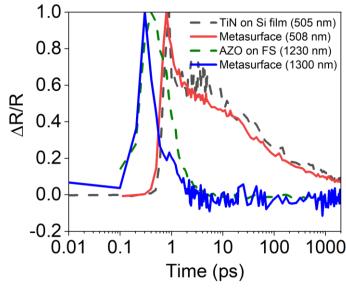
- Permittivity Tuning with increased pump fluence
- Large modulation in the NIR

TiN-AZO

- Fast and Slow switching in the same Berreman metasurface
- <10 ps to >3 ns switching time by controlling the probe







Acknowledgement

Work on Y:CdO published in Advanced Functional Materials



Work on ZnO modulators preprint in arxiv





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